

# Examination of Current Profiles in Magnetic Islands During RF Current Condensation

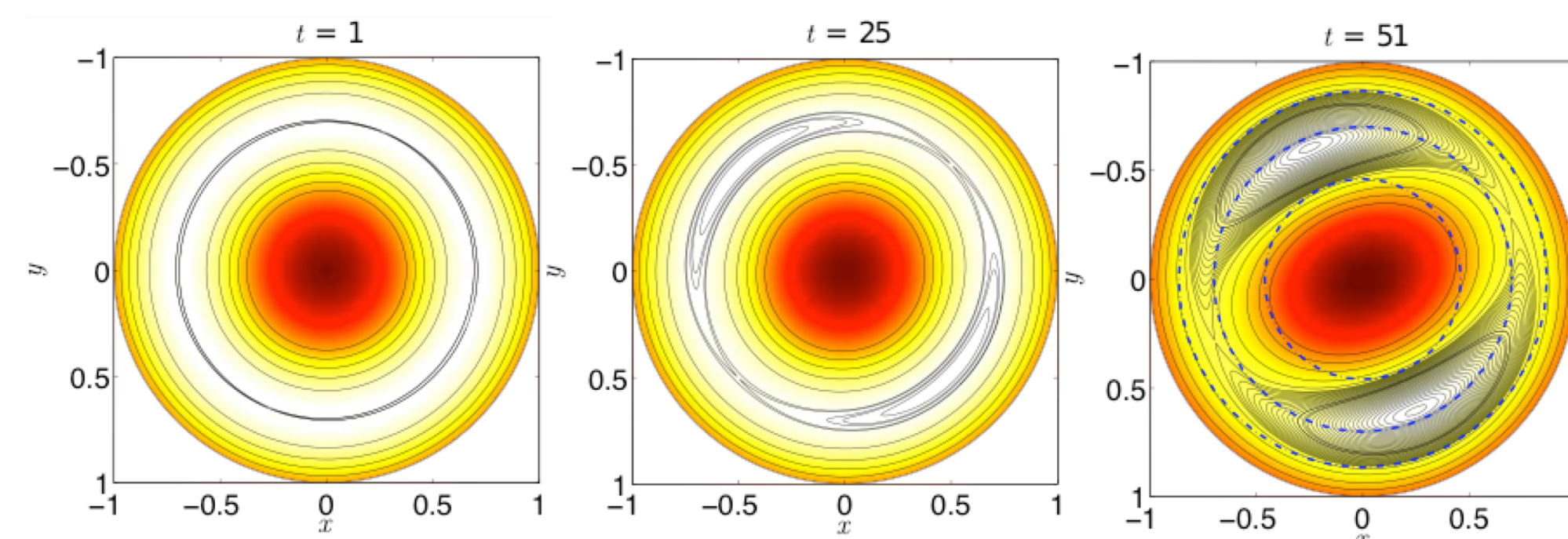
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## Magnetic islands degrade confinement in fusion devices

- Magnetic islands are formed by the tearing of toroidally nested flux surfaces

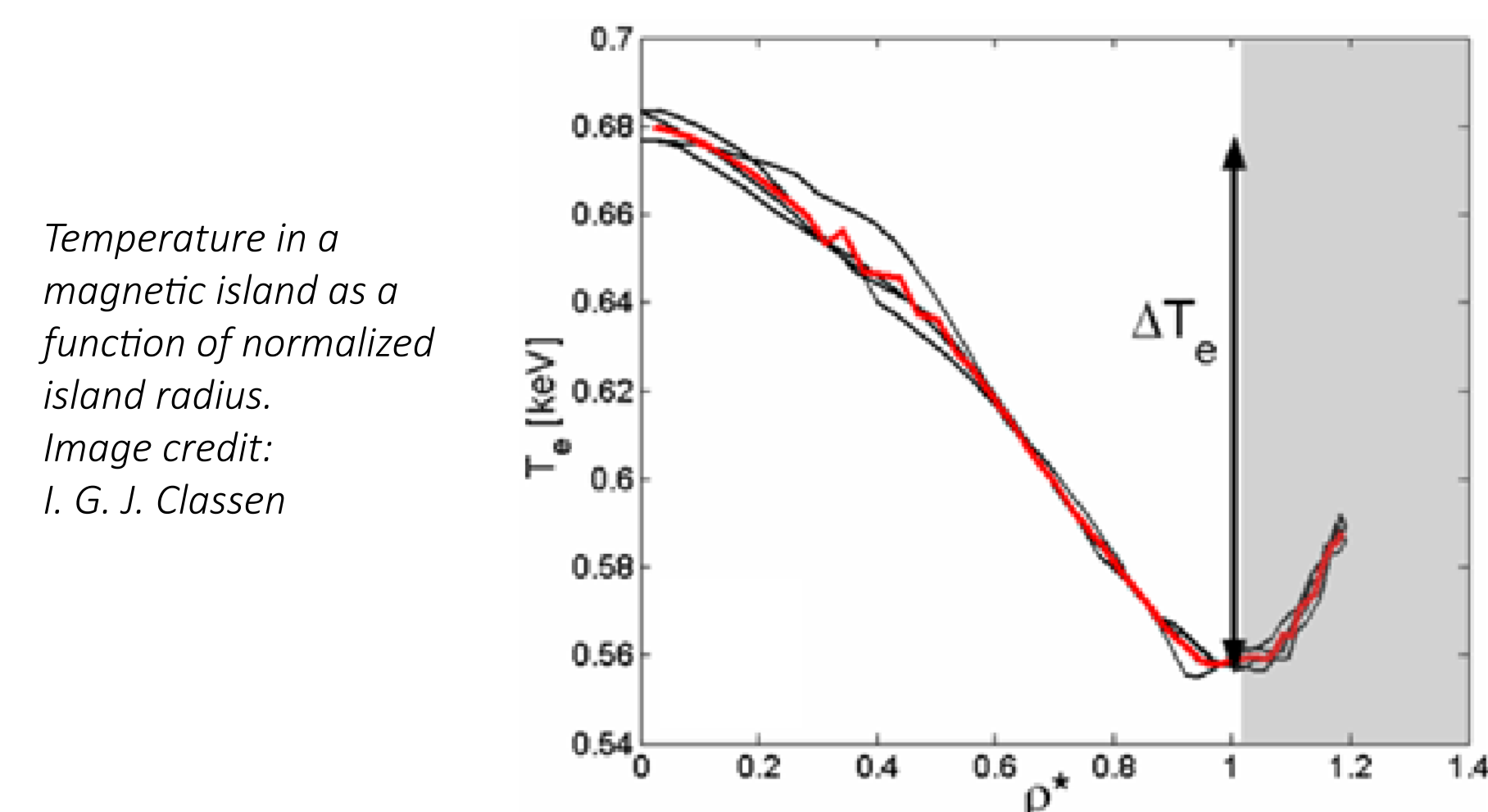


Formation of magnetic islands in a tokamak.  
Image credit: Qian Teng

- Magnetic islands are thought to be destabilized by a loss of bootstrap current
- Islands can be stabilized by driving current directly into their centers
- Typically, electron cyclotron current drive (ECCD) and lower hybrid current drive (LHCD) are used
- However, radio frequency (RF) driven current needs to be constantly readjusted to accommodate island formation and movement

## The RF current condensation effect could stabilize magnetic islands without the need for precise steering of waves

- Both electron cyclotron and lower hybrid waves deposit the most power where electron temperature is greatest
- In a magnetic island, electron temperature is greatest at the island center

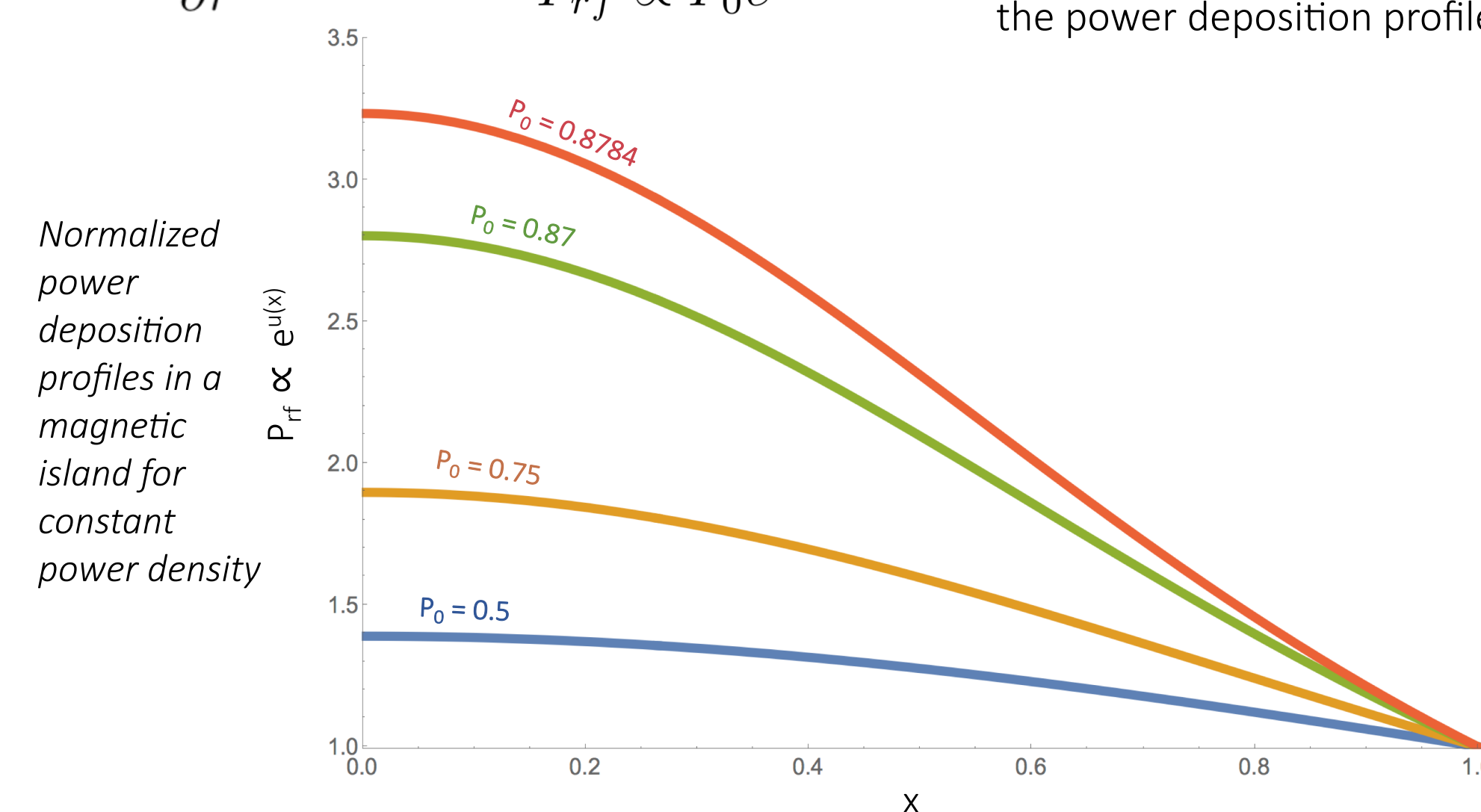


Temperature in a magnetic island as a function of normalized island radius.  
Image credit: I. G. J. Classen

- As a result, ECCD and LHCD should naturally deposit more power at the island center
- This should cause electron temperature to increase further, resulting in a positive feedback loop and the RF current condensation effect

## Constant power density results in noticeable current condensation at the island center

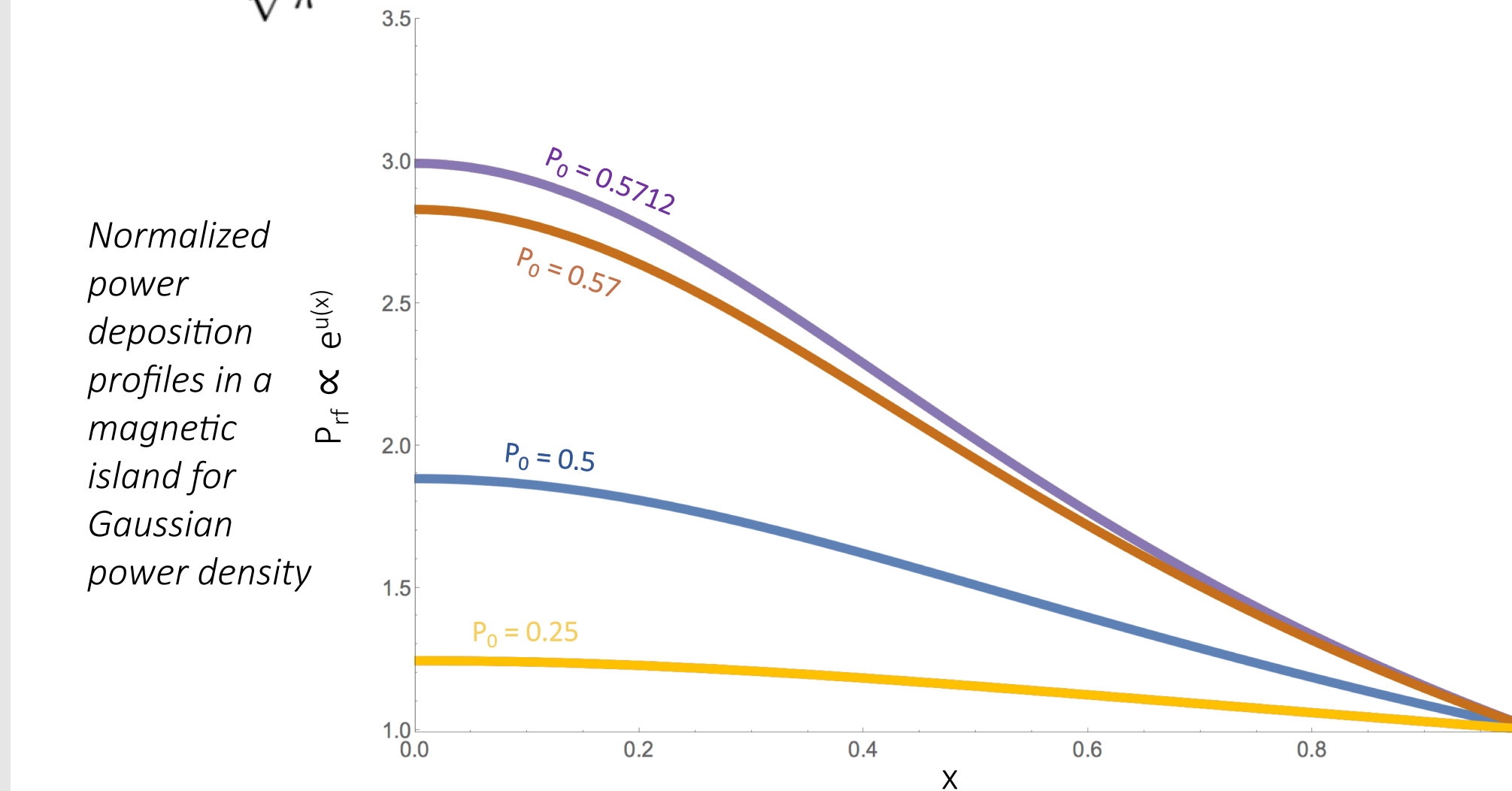
- Assume that the power density available for absorption by the island is constant
- Steady-state diffusion equation in 1D slab island:  $n\kappa_{\perp} \frac{\partial^2 T}{\partial r^2} = -P_{rf}$
- Power dissipated in the island for constant power density:  $P_{rf} \propto P_0 e^{\omega_{1s}^2 \bar{T}/T_s}$
- Setting  $u \equiv \omega_{1s}^2 \bar{T}/T_s$  and normalizing other variables, can solve the diffusion equation for  $u(x)$  to obtain the power deposition profiles



- Constant power density demonstrates notable current condensation at the island center

## Gaussian power density causes similar power deposition at center as constant power density case

- Assume that the power density takes the form of a Gaussian, centered at the island center
- Power dissipated for Gaussian power density of  $e^{-4x^2}$ :  $P_{rf} \propto \frac{4}{\sqrt{\pi}} P_0 e^{-4x^2} e^{\omega_{1s}^2 \bar{T}/T_s}$
- Employing the same diffusion equation and process as the constant case, can obtain  $u(x)$  and the deposition profiles



- Gaussian power density also exhibits significant current condensation at the island center

## RF current condensation holds potential for stabilization of islands

- All three scenarios examined demonstrate noticeable current condensation at the island center
- Wave depletion for large  $V_0$  recovers the constant power density deposition profiles
- Wave depletion for small  $V_0$  can allow for increased current condensation
- Wave depletion above the bifurcation point could result in even greater current condensation at the island center

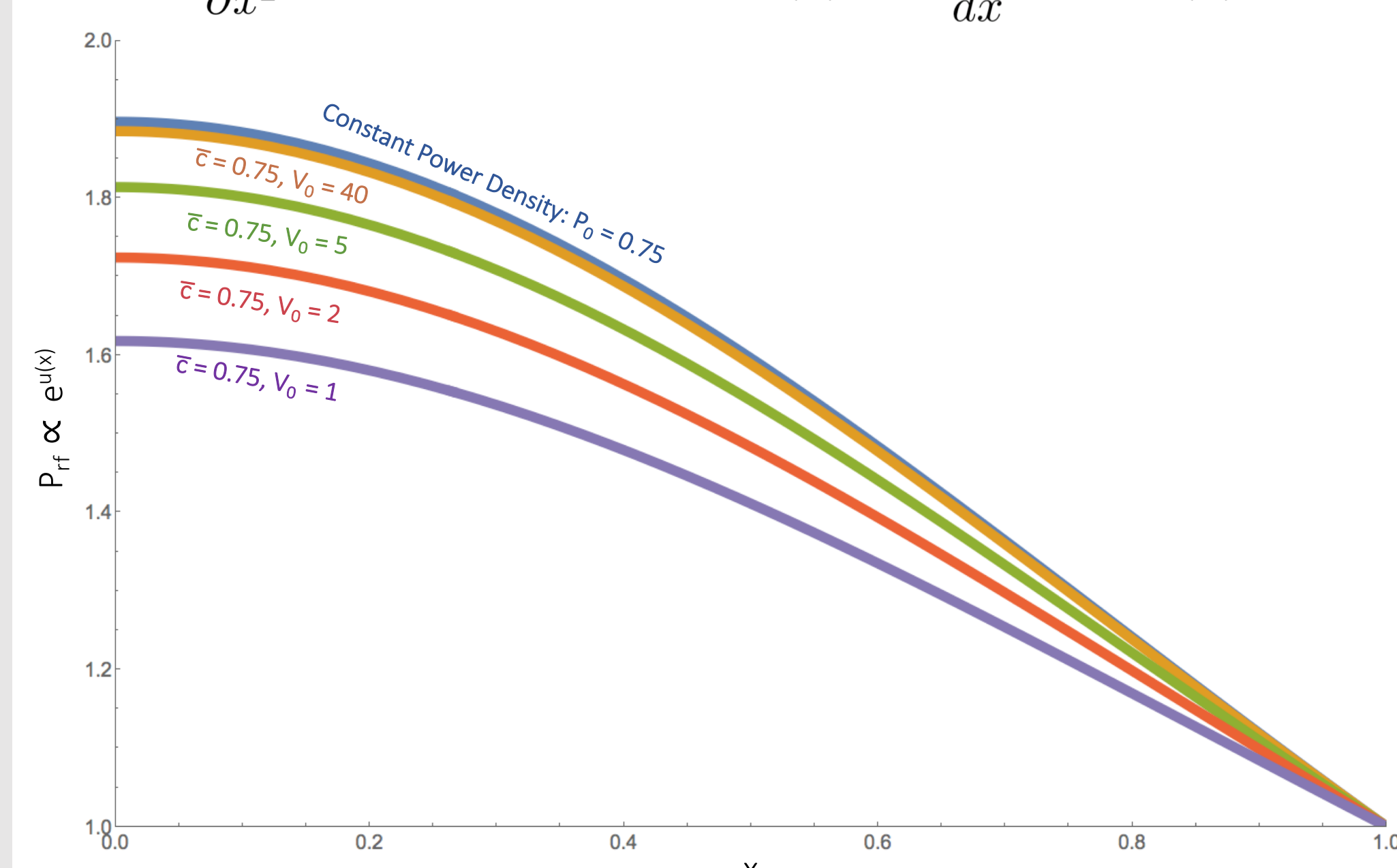
## Future work

- Further explore the impact of  $V_0$  on the bifurcation threshold
- Explore deposition profiles above the bifurcation threshold
- Construct profiles including the symmetrizing term,  $V'(-x)$
- Experimentally verify the existence of the RF current condensation effect in a tokamak

## Accounting for wave depletion can result in shifting of the bifurcation point and heightened current condensation at the island center

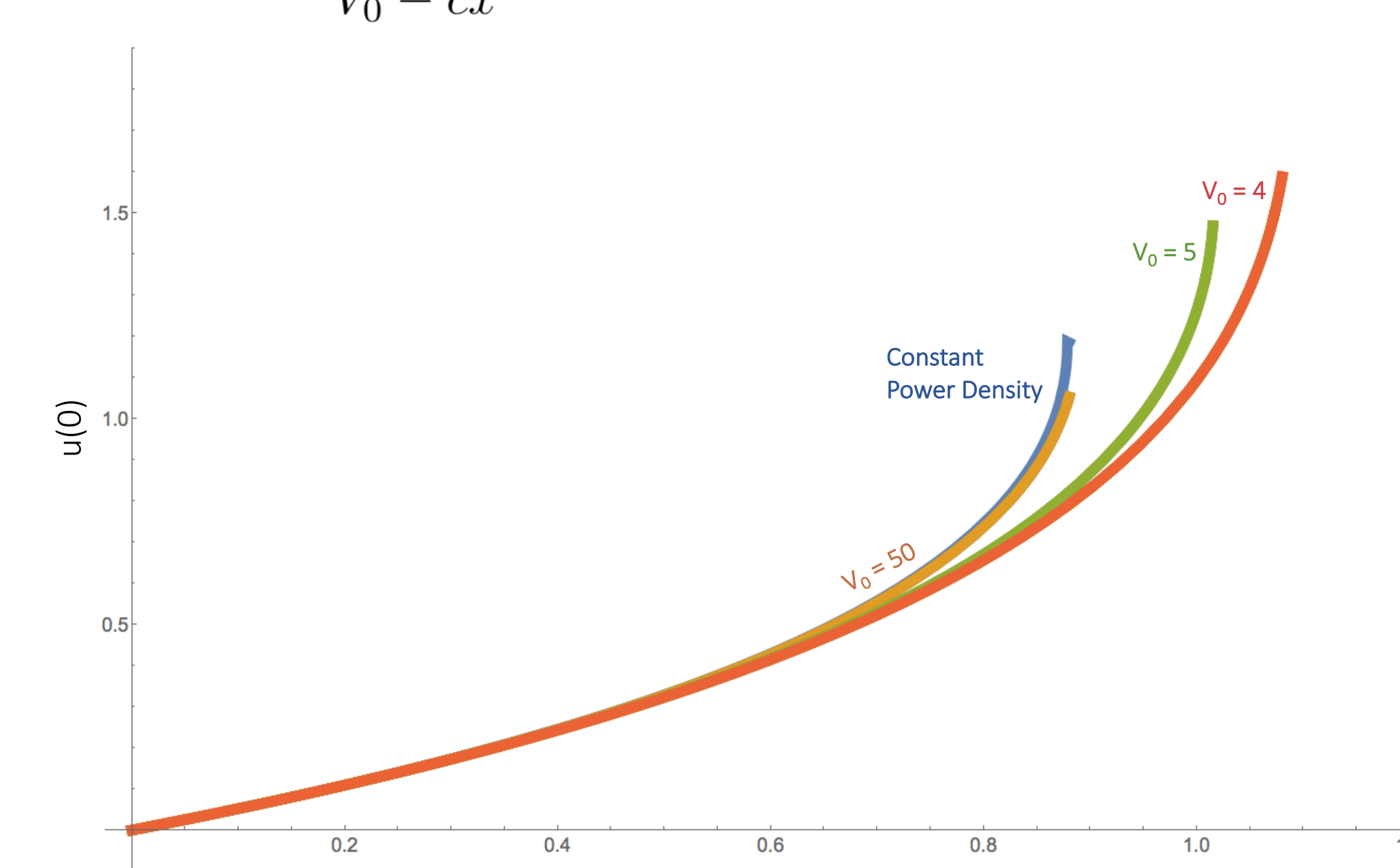
- Wave depletion model accounts for the energy lost by the RF waves as they travel through the island
- Results in a more realistic treatment of the power deposition than the above approaches

- Steady-state diffusion equation, with wave depletion:  $n\kappa_{\perp} \frac{\partial^2 T}{\partial x^2} = \bar{V}'(x) + \bar{V}'(-x)$
- Change in energy density of the RF wave:  $\bar{V}'(x) = \frac{d\bar{V}(x)}{dx} = -\alpha(x) e^{\omega_{1s}^2 \bar{T}/T_s} \bar{V}$
- Spatial damping of the wave in the linear limit:  $\alpha(x) = \frac{\bar{c}}{V_0 - \bar{c}x}$
- Setting  $u \equiv \omega_{1s}^2 \bar{T}/T_s$ ,  $V \equiv \omega_{1s}^2 \bar{V}/(n\kappa_{\perp} T_s)$ , and neglecting  $V'(-x)$ , can solve the diffusion equation for  $u(x)$  to obtain the power deposition profiles



Normalized power deposition profiles for both constant power density and wave depletion models

- As  $V_0$  increases, wave depletion profiles approach constant power density profiles
- For small  $V_0$ , wave depletion shifts the bifurcation point and can allow for greater current condensation at the island center



$u(0)$  as a function of either  $\bar{c}$  or  $P_0$ , where  $u(0)$  is the value of  $u$  at the island center. The curves end at their respective bifurcation points

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## References

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